# **Quantum Mechanics**

# **Identical particles**

**Problem 1.-** Find the expectation value  $\langle (x_1 - x_2)^2 \rangle$  for two particles in the simple harmonic oscillator potential in states  $|0\rangle$  and  $|1\rangle$  in the cases

- a) If the particles are distinguishable.
- b) If they are identical bosons.
- c) If they are identical fermions.

## Solution:

Let's call  $|n\rangle$  the eigen function of the harmonic oscillator in one dimension with energy eigen value  $(n+1/2)\hbar\omega$ . The quantum number n is a non-negative integer.

To calculate the various expectation values, we will need:

$$x = \sqrt{\frac{\hbar}{2m\omega}} (a^+ + a^-)$$
 This operator is obtained by adding  $a^+ + a^-$ 

To apply the operator, recall that:

$$a^{+}|n\rangle = \sqrt{n+1}|n+1\rangle$$
$$a^{-}|n\rangle = \sqrt{n}|n-1\rangle$$

The expectation value of the square of x can be found using this operator:

$$\langle x^2 \rangle = \langle n | x^2 | n \rangle = \langle n | \left( \sqrt{\frac{\hbar}{2m\omega}} (a^+ + a^-) \right)^2 | n \rangle = \frac{\hbar}{2m\omega} \langle n | (a^+ + a^-)^2 | n \rangle$$

$$\langle x^2 \rangle = \frac{\hbar}{2m\omega} \langle n | (a^{+2} + a^- a^+ + a^+ a^- + a^{-2}) | n \rangle =$$

$$= \frac{\hbar}{2m\omega} \langle n | (\sqrt{(n+1)(n+2)} | n+2 \rangle + (2n+1) | n \rangle + \sqrt{n(n-1)} | n-2 \rangle )$$

Notice that only the term in the middle gives us a nonzero result, which is equal to:

$$\left\langle x^2 \right\rangle = \frac{\hbar}{2m\omega}(2n+1)$$

In our problem:  $\langle 0|x^2|0\rangle = \frac{\hbar}{2m\omega}$  and  $\langle 1|x^2|1\rangle = \frac{3\hbar}{2m\omega}$ 

We also know, just from symmetry that the expectation value of x vanishes for a harmonic oscillator.

To get the "cross integral"  $\langle 0|x|1 \rangle$  we again use the operator above:

$$\langle 0|x|1\rangle = \langle 0|\left(\sqrt{\frac{\hbar}{2m\omega}}(a^+ + a^-)\right)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(a^+ + a^-)|1\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle) = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2\rangle + 1|0\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle 0|(\sqrt{2}|2$$

#### a) Distinguishable particles:

$$\left\langle (x_1 - x_2)^2 \right\rangle = \left\langle x^2 \right\rangle_a + \left\langle x^2 \right\rangle_b - 2\left\langle x \right\rangle_a \left\langle x \right\rangle_b = \frac{\hbar}{2m\omega}(2(0) + 1) + \frac{\hbar}{2m\omega}(2(1) + 1) - 0 = \frac{2\hbar}{m\omega}$$

## b) Identical Bosons:

$$\left\langle (x_1 - x_2)^2 \right\rangle = \left\langle x^2 \right\rangle_a + \left\langle x^2 \right\rangle_b - 2\left\langle x \right\rangle_a \left\langle x \right\rangle_b - 2\left| \left\langle x \right\rangle_{ab} \right|^2 = \frac{\hbar}{2m\omega} (2(0) + 1) + \frac{\hbar}{2m\omega} (2(1) + 1) - 0 - 2\frac{\hbar}{2m\omega} = \frac{\hbar}{m\omega}$$

#### b) Identical Fermions in the triplet state:

$$\left\langle (x_1 - x_2)^2 \right\rangle = \left\langle x^2 \right\rangle_a + \left\langle x^2 \right\rangle_b - 2\left\langle x \right\rangle_a \left\langle x \right\rangle_b + 2\left| \left\langle x \right\rangle_{ab} \right|^2 = \frac{\hbar}{2m\omega} (2(0) + 1) + \frac{\hbar}{2m\omega} (2(1) + 1) - 0 + 2\frac{\hbar}{2m\omega} = \frac{3\hbar}{m\omega}$$

**Problem 2.-** Consider that two particles interact through a potential that only depends on the vector  $\vec{r} = \vec{r_1} - \vec{r_2}$  and show that the Schrodinger equation can be separated in two equations, one for  $\vec{r}$  and the other for the center of mass vector  $\vec{R} = \frac{m_1 \vec{r_1} + m_2 \vec{r_2}}{m_1 + m_2}$ 

**Solution:** We write the kinetic energy in terms of the reduced mass,  $\mu = \frac{m_1 m_2}{m_1 + m_2}$ , the total mass and the vectors  $\vec{R}$  and  $\vec{r}$ :

$$-\frac{\hbar^2}{2(m_1+m_2)}\nabla_R^2\Psi - \frac{\hbar^2}{2\mu}\nabla_r^2\Psi + V(r)\Psi = E\Psi$$

Where  $\Psi = \psi_R(R)\psi_r(r)$ 

$$-\frac{\hbar^2}{2(m_1+m_2)}\nabla_R^2\psi_R(R)\psi_r(r) - \frac{\hbar^2}{2\mu}\nabla_r^2\psi_R(R)\psi_r(r) + V(r)\psi_R(R)\psi_r(r) = E\psi_R(R)\psi_r(r)$$
$$-\frac{\hbar^2}{2(m_1+m_2)}\psi_r(r)\nabla_R^2\psi_R(R) - \frac{\hbar^2}{2\mu}\psi_R(R)\nabla_r^2\psi_r(r) + V(r)\psi_R(R)\psi_r(r) = E\psi_R(R)\psi_r(r)$$

Dividing by  $\Psi = \psi_R(R)\psi_r(r)$ :

$$-\frac{\hbar^2}{2(m_1+m_2)}\frac{\nabla_R^2\psi_R(R)}{\psi_R(R)} - \frac{\hbar^2}{2\mu}\frac{\nabla_r^2\psi_r(r)}{\psi_r(r)} + V(r) = E$$
$$-\frac{\hbar^2}{2(m_1+m_2)}\frac{\nabla_R^2\psi_R(R)}{\psi_R(R)} = E_R$$
$$-\frac{\hbar^2}{2\mu}\frac{\nabla_r^2\psi_r(r)}{\psi_r(r)} + V(r) = E_r$$

and  $E = E_r + E_R$ 

The two separate equations are:

$$-\frac{\hbar^2}{2(m_1+m_2)}\nabla_R^2\psi_R(R) = E_R\psi_R(R)$$
$$-\frac{\hbar^2}{2\mu}\nabla_r^2\psi_r(r) + V(r)\psi_r(r) = E_r\psi_r(r)$$

and  $E = E_r + E_R$